

# Testing Compliance of WPT Devices with DASY8/6 Module SAR and Module WPT according to IEC/IEEE 63184

#### 1 Scope of this Document

This application note provides guidance on how to demonstrate compliance of inductive wireless power transfer (WPT) systems in the laboratory with dimensions of not larger than  $900 \times 900 \,\mathrm{mm^2}$  with basic restrictions (BR; Tier 4) and reference levels (RL; Tier 2) in accordance with the latest draft of IEC/IEEE 63184 [1]<sup>1</sup>. It represents the most accurate assessment methods currently available, as defined in Sections 6 and 8 of [1]:

- ≥4 MHz: DASY8/6 Module SAR 16.4+ measures and determines the induced electric (E-)field and specific absorption rate (SAR) inside the standard phantom directly and fully complies with Section 6.2 of [1];
- <4 MHz: DASY8/6 Module WPT V2.6+ (i) measures the incident fields, including phase, in a volume equivalent to the entire half-space at high resolutions, (ii) reconstructs the Maxwell field in this volume generated by the complete WPT system that may include the transmitter and receiver, e.g., the wireless charger with the phone on top, and (iii) reconstructs the induced SAR and E-fields inside the standard phantom by means of the fully validated Sim4Life solver (P-EM-QS). The workflow does not require any modeling of the WPT system by the user. The measurement system is self-contained, i.e., it automatically verifies if all conditions for reliable and accurate evaluations are satisfied, and provides the analysis of compliance according to the latest regulations. It fully complies with the requirements of Sections 6 and 8 of IEC/IEEE 63184 [1].

# 2 Summary of IEC/IEEE 63184

#### 2.1 Scope and Method

IEC/IEEE 63184 specifies the methods to demonstrate compliance of inductive WPT systems working in the frequency range from 3 kHz to 30 MHz against BR and RL. Regarding the assessment of the direct effects (induced E-field, current density, SAR), IEC/IEEE 63184 provides a 4-Tier approach, as shown in Figure 1.1. Tier 1 generally greatly overestimates the exposure while Tier 4 is the most accurate evaluation with the least overestimation. Tier 1 and 2 compare the incident fields generated by magnetic near-field sources with the RL. Tiers 3 and 4 determine the compliance directly against the BR, i.e., peak induced E-fields ( $pE_{ind}$ ) and peak spatially averaged SAR over a mass of 1 g or 10 g (psSAR1g/10g). Tier 3 estimates the induced fields with a coupling factor based on gradients of the incident magnetic (H-)field, e.g., [3], whereas Tier 4 requires the induced fields to be determined directly using computational or experimental techniques.

In Section 6 of IEC/IEEE 63184, requirements for measurement methods are provided. For incident field measurements, three-axis sensors are recommended. To improve the measurement accuracy at short distances, the sensors need to be sufficiently small. H-field sensors with a loop size of  $\leq 1\,\mathrm{cm}^2$  is recommended. The MAGPy V2.0 probe used by DASY8/6 Module WPT V2.6+ meets these requirements. For SAR measurements, IEC/IEEE

 $<sup>^1</sup>$ Other procedures will be available for evaluating other systems such as capacitive or radiative WPT systems.

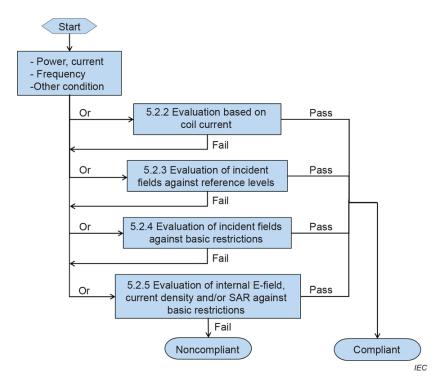


Figure 1.1: Flowchart of the assessment procedure to demonstrate compliance of inductive WPT systems specified in IEC/IEEE 63184 [1].

63184 refers to the procedures specified in IEC/IEEE 62209-1528. DASY8/6 Module SAR 16.4+ enables direct SAR measurements and is fully compatible with IEC/IEEE 62209-1528. IEC/IEEE 63184 specifies the dielectric properties of the tissue simulating liquid for SAR measurements. The homogeneous phantom used by DASY8/6 Module WPT V2.6+ follows this medium definition.

In Section 8, IEC/IEEE 63184 defines the requirements for hybrid methods which combine measurement and computational methods. These methods do not require any device under test (DUT) model, and hence remove the associated large uncertainties. DASY8/6 Module WPT V2.6+ is fully compatible with the requirements of Section 8 of [1].

#### 3 SPEAG's Measurement Solutions

#### 3.1 DASY8/6 Module SAR 16.4+

DASY8/6 Module SAR 16.4+ meets all performance requirements of IEC/IEEE 62209-1528:2020 [2] for frequencies between 4 MHz and 10 GHz. More details about DASY8/6 Module SAR 16.4+ are provided in the DASY8/6 Module SAR 16.4+ Manual [8]. DASY8/6 Module SAR 16.4+ allows direct measurements of SAR and local induced E-field for compliance testing against BR, so it corresponds to the Tier 4 procedure of IEC/IEEE 63184.

#### 3.2 DASY8/6 Module WPT V2.6+

DASY8/6 Module WPT V2.6+ meets all performance requirements of IEC/IEEE 63814 [1]. It is composed of the isotropic probe MAGPy-8H3D+E3D Version 2, the reference probe (MAGPy-RA $\phi$ V2), and the data acquisition system (MAGPy-DAS) mounted to the DASY8/6 robot via the emergency stop (MAGPy-ES). At each probe location, eight isotropic H-field values plus the phase are acquired in addition to the isotropic E-field measurement.

DASY8/6 Module WPT V2.6+ offers the Tier 2 procedure of IEC/IEEE 63184 with the measurement method. The incident H-field is measured on a high density grid (7.33 mm resolution) such that the incident quasi-static H-field (amplitude and phase) in the entire measured volume can be reconstructed by means of our advanced and validated vector potential reconstruction (see Appendix A for more information). The incident E-field distribution is measured in the same volume, enabling accurate determination of the field impedance at  $d=30 \, \text{mm}$ . Due to the geometric design of the Ez-field sensor, the measured information is sufficient for a reliable estimation of the E-fields at the surface of the DUT, i.e., the distance d=0, and its potential coupling to the tissue simulating media, even for very localized E-field sources. The effect of the phantom loading or backscattering is less than 1% for frequencies <4 MHz as derived from theoretical considerations, supported by simulations and verified by measurements (see Appendix B).

DASY8/6 Module WPT V2.6+ also offers the Tier 4 procedure of IEC/IEEE 63184 according to Section 8. The measured and reconstructed fields are used to assess the induced fields due to the incident H-field, without approximation and with known uncertainty, by Sim4Life's Quasi-Static EM Solver (P-EM-QS) (ZMT Zurich MedTech AG, Zurich, Switzerland). Since neglecting the effect of the phantom loading or backscattering results in a small overestimation of the induced fields, the assessment with DASY8/6 Module WPT V2.6+ is conservative. The induced fields due to the incident E-field² is determined by a conservative approximation that is valid for local E-field sources [9]. The validity of the local E-field condition is automatically assessed by the system, including checking whether the field impedance is less than 10% of the plane wave impedance of 377  $\Omega$ . The total field evaluation (see Appendix C for its validation) provides the assessed total induced fields, which are compared to BR.

The dedicated graphical user interface (GUI) fully automates the testing workflow. More details about DASY8/6 Module WPT V2.6+ are provided below and in the DASY8/6 Module WPT V2.6+ Manual [10].

 $<sup>^2</sup>$ The strongest E-field generated by a WPT system is often traceable to local accumulation of charge, e.g., across a discrete capacitor (to achieve resonance) and at the end of conductors, that decays rapidly as a function of d at a rate of  $1/d^4$  but can potentially induce fields in the body [9]. The problem that these charge accumulations are difficult to predict or accurately simulate is overcome with DASY8/6 Module WPT V2.6+, which determines the field characteristics with measurements.

# 4 Test System and Procedures for Frequencies ≥4 MHz

#### 4.1 System Requirements

To determine the induced instant peak E-field ( $pE_{ind,inst}$ , as specified in [12]) and the psSAR values (psSAR1g/10g, as specified in [13]), the following system configuration is recommended:

- DASY8/6 Module SAR software
- ELI phantom
- HBBL4-250Vx head simulating liquid
- EX3DVx probe with conversion factor assessment at 6 MHz (covers 4–9 MHz) and 13 MHz (covers 9–19 MHz)
- confined loop antennas CLA-6 and CLA-13 for system check and validation purposes

#### 4.2 Measurement Procedure

The workflow to demonstrate compliance with the BR, illustrated in Figure 1.2, is equivalent to standard psSAR evaluations with an additional step – Determination of  $pE_{ind,tissue,inst}$  – described below.

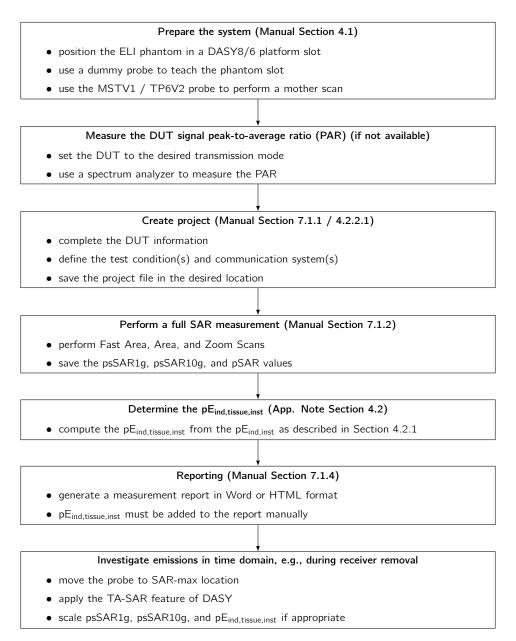


Figure 1.2: Step-by-step measurement procedure for using DASY8/6 Module SAR to evaluate compliance of WPT devices and systems with BR.

#### 4.2.1 Determining pE<sub>ind,tissue,inst</sub>

The psSAR1g/10g is determined from a regular Zoom Scan. In addition, DASY8/6 Module SAR 16.4+ reports the pSAR (maximum SAR at the inner phantom surface) and the pE $_{ind,inst}$  (see Figure 1.3).



Figure 1.3: pE<sub>ind,tissue,inst</sub> results derived from a Zoom Scan measured with DASY8/6 Module SAR.

The maximum instantaneous E-field value in non-homogeneous tissues can be calculated using Equation 1.1,

$$pE_{ind,tissue,inst} = pE_{ind,inst} \cdot CF \tag{1.1}$$
 where 
$$pE_{ind,inst} = pE_{ind,inst} \cdot CF \tag{1.1}$$
 where 
$$pE_{ind,inst} = maximum instantaneous E-field value at the phantom surface in V/m = coverage factor, which translates the induced fields in the standardized homogeneous tissue to those in non-homogeneous tissues, after taking field enhancements at dielectric boundaries into account, as generally determined in [4] 
$$pE_{ind,tissue,inst} = maximum instantaneous E-field values in non-homogeneous tissues$$$$

**Note:** The 5 mm-line and  $2\times2\times2$  mm<sup>3</sup>-cube averaged induced E-field will be directly provided in DASY8/6 Module SAR V17.0.

#### 4.3 Uncertainty

The uncertainty for evaluations performed with DASY8/6 Module SAR 16.4+ was determined according to IEC/IEEE 62209-1528:2020 [2] and documented in the DASY8/6 Module SAR 16.4+ Manual [8]. Typically, the uncertainties (k=2) are <22.8% for psSAR1g/10g (see Table 1.1) and <12.5% for pE<sub>ind,inst</sub>, based on the assumption of a maximum extrapolation uncertainty of less than 5%.

# DASY8/6 Uncertainty Budget According to IEC/IEEE 62209-1528

(Frequency band: 4 MHz–300 MHz)

		Unc.	Prob.	Div.	$(c_i)$	$(c_i)$	Std. Unc.	Std. Unc.
Symbol	Error Description	Value	Dist.		(1g)	(10 g)	(1g)	(10 g)
	nent System Errors							
CF	Probe Calibration	±13.3%	N	2	1	1	±6.65%	±6.65%
CF <sub>drift</sub>	Probe Calibration Drift	±1.7%	R	$\sqrt{3}$	1	1	±1.0%	±1.0%
LIN	Probe Linearity	±4.7%	R	$\sqrt{3}$	1	1	±2.7%	±2.7%
BBS	Broadband Signal	±0.6%	R	$\sqrt{3}$	1	1	±0.3%	±0.3%
ISO	Probe Isotropy	±7.6%	R	$\sqrt{3}$	1	1	±4.4%	±4.4%
DAE	Other Probe+Electronic	±0.8%	N	1	1	1	±0.8%	±0.8%
AMB	RF Ambient	±1.8%	N	1	1	1	±1.8%	±1.8%
$\Delta_{sys}$	Probe Positioning	±0.006 mm	N	1	0.04	0.04	±0.10%	±0.10%
DAT	Data Processing	±1.2%	N	1	1	1	±1.2%	±1.2%
Phantom	and Device Errors							
$LIQ(\sigma)$	Conductivity (meas.)	±2.5%	N	1	0.78	0.71	±2.0%	±1.8%
$LIQ(T_{\sigma})$	Conductivity (temp.)	±5.4%	R	$\sqrt{3}$	0.78	0.71	±2.4%	±2.2%
EPS	Phantom Permittivity	$\pm 14.0\%$	R	$\sqrt{3}$	0	0	±0%	±0%
DIS	Distance DUT – TSL	±2.0%	N	1	2	2	±4.0%	±4.0%
$D_{xyz}$	Device Positioning	±1.0%	N	1	1	1	±1.0%	±1.0%
Н	Device Holder	±3.6%	N	1	1	1	±3.6%	±3.6%
MOD	DUT Modulation	±2.4%	R	$\sqrt{3}$	1	1	±1.4%	±1.4%
TAS	Time-average SAR	$\pm 1.7\%$	R	$\sqrt{3}$	1	1	±1.0%	±1.0%
RF <sub>drift</sub>	DUT drift	±2.5%	N	1	1	1	±2.5%	±2.5%
VAL	Val Antenna Unc.	±0.0%	N	1	1	1	±0%	±0%
RFin	Unc. Input Power	±0.0%	N	1	1	1	±0%	±0%
Correction	to the SAR results							
$C(\varepsilon, \sigma)$	Deviation to Target	±1.9%	N	1	1	0.84	±1.9%	±1.6%
C(R)	SAR scaling	±0.0%	R	$\sqrt{3}$	1	1	±0.0%	±0.0%
u(ΔSAR)	Combined Uncertainty						±11.4%	±11.3%
U	Expanded Uncertainty						±22.8%	±22.5%

Table 1.1: Uncertainty budget for peak  $1\,g$  and  $10\,g$  mass-averaged SAR measured with DASY8/6 Module SAR, assessed according to IEC/IEEE 62209-1528.

#### 5 Test System and Procedures for Frequencies <4 MHz

#### 5.1 System Requirements

To determine induced fields for compliance testing against BR or to determine incident fields for coompliance testing against RL, the following equipment is required:

- DASY8/6 Module WPT V2.6+ including:
  - MAGPy-8H3D+E3D Version 2 probe with the integrated data acquisition system MAGPy-DAS
  - MAGy-RA $\phi$ V2 reference probe as a phase reference
  - MAGPy-ES emergency stop system
- WPT sources (incl. V-Coil500/3, V-Coil350/85, and V-Coil50/400) for system check and validation purposes
- Software DASY6/8 Module WPT V2.6+

#### 5.2 Assessment of Peak Induced Fields

#### 5.2.1 Measurement Procedure

The workflow to demonstrate compliance with the BR is illustrated in Figure 1.4. Detailed descriptions of each step can be found in Section 7 of the DASY8/6 Module WPT Manual [10]. It is recommended to perform a system check before any compliance testing with the V-Coil source that operates at the frequencies closest to that of the DUT. This provides the confidence that the system operates within its specifications.

# Prepare system (Manual Sections 7.1.1 and 7.1.2) • install the MAGPy probe and align it in the light beam connect and position the reference probe and verify that the signal-to-noise ratio is sufficient Inspect signal (Manual Section 7.1.2) • move the probe tip to DUT surface check the waveforms and spectra of incident H- and E-fields check default parameters of time-domain slicing and peak frequency searching and adjust if needed Configure volume scan (Manual Section 7.1.2) • teach DUT position • specify extent of scan or enable auto-extension mode Measure incident fields by volume scan (Manual Section 7.1.2) • start 3D scan of H- and E-fields (amplitude, frequency, and phase evaluation at each • verify the stability of the signal Determine induced fields (Manual Section 7.1.2) • start post processing: 1) determine incident H-field at the probe tip by surface field reconstruction; 2) determine induced fields by vector potential reconstruction and simulation in a standardized phantom • perform a sanity check of the reconstruction and simulation Evaluate compliance of H-field source only (Manual Section 7.1.3, see Figure 1.5) • result tables showing compliance against limits from different safety standards option to apply the multi-frequency enhancement factor and remove undesired peak frequencies (e.g., those confirmed to be noise) · option to apply the coverage factor Evaluate compliance for total field (Manual Section 7.1.4, see Figure 1.7) • verify that the incident E-field is local (see Figure 1.6) • option to include the induced fields due to the incident E-field result tables show compliance of total induced fields against limits from different safety standards Reporting (see Figure 1.8)

Figure 1.4: Step-by-step measurement procedure for using DASY8/6 Module WPT to evaluate compliance of WPT devices and systems with BR.

• check the field amplitude in time domain and scale results in the report if appropriate

Investigate emissions in time domain, e.g., during receiver removal (see Figure 1.9)

• move the MAGPy probe to H-max location

generate a report

open MAPGy software

Simulation			Res	ults			Complia	ance (Fie	ld value	Com	plianc	e (Ratio	s)	Frequ	iency-do	main	Ti	me-dor	nain							
■ Total fi	eld eva	luation	1			<b>✓</b> A	pply cov	erage fac	tor				Multi-	-frequen	ıcy enha	ncement			■ Disp	olay ratio	s in dB					
			ICNIRE	2010/20	020			ICNIR	P 1998				IEI	EE 2019					cc				н	Code 6		
			RL			BR		RL		BR		E	RL			DRL	N	IPE		BR			RL			BR
Distance [mm]	pl	Hinc	F	Einc	pE <sub>ind</sub>	psSAR	pH <sub>inc</sub>	pE <sub>inc</sub>	pJ <sub>ind</sub>	psSAR	Р	H <sub>inc</sub>	P	Einc	pEind	psSAR	pH <sub>inc</sub>	pE <sub>inc</sub>	pE <sub>ind</sub>	psSAR	P	H <sub>inc</sub>		pE <sub>inc</sub>	pE <sub>ind</sub>	psSAR
	NS	тн	NS	тн	NS	тн	N/A	N/A	NS	тн	NS	тн	NS	тн	NS	тн	N/A	N/A	N/A	тн	NS	тн	NS	тн	NS	TH
2.00	11.8	20.2	88.1	38.5	0.33	<0.01	135.0	290.0	2.1	<0.01	1.52	2.72	11.9	90.0	0.12	<0.01	152.0	70.7		<0.01	2.75	135.0	88.1	433.0	0.50	<0.01
									v	Eind, cube a	<sub>g.</sub> = [5.	.34], w <sub>Eir</sub>	nd, local =	= [7.54],	WE <sub>ind, line</sub>	<sub>avg.</sub> = [2.9										

Figure 1.5: Table in DASY8/6 Module WPT showing the compliance evaluation results for the induced SAR due to the incident H-field without condering any potential E-field source (note that the "Total field evaluation" option (highlighted by the orange box) is un-checked here).

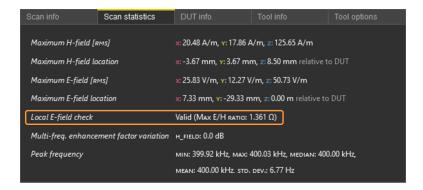


Figure 1.6: The section of the DASY8/6 Module WPT GUI that displays the statistical information of the volume scan, e.g., the result of the local E-field check (highlighted by the orange box; the maximum E/H ratio is also listed).



Figure 1.7: Table in the DASY8/6 Module WPT GUI showing the compliance evaluation results for the total induced SAR (i.e., for exposures from both incident H- and E-fields; note that the "Total field evaluation" option (highlighted by the orange box) is checked here).

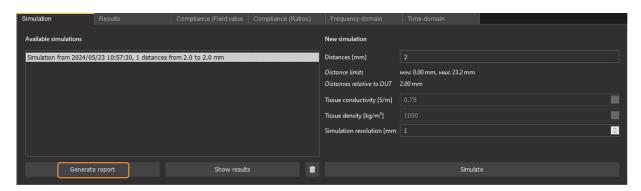


Figure 1.8: The section of the DASY8/6 Module WPT GUI for setting up the simulation and generating the report. Note that the tissue conductivity and mass density are pre-set to standardized values and cannot be altered by users. After the simulation is completed, a report can be generated by clicking the "Generate report" button (highlighted by the orange box).



Figure 1.9: The time-domain plot of the incident H-field in the MAGPy GUI. The data were recorded from a commercial wireless charger while removing the smartphone placed on the charger. The same procedure can also be used to monitor the stability of the source.

#### 5.2.2 Uncertainty

The uncertainty for evaluations performed with DASY8/6 Module WPT V2.6+ was determined according to IEC/IEEE 63184 [1] and documented in the DASY8/6 Module WPT V2.6+ Manual [10]. Typically, the uncertainty (k=2) is <33.9% for psSAR1g/10g (see Tables 1.2 and 1.3) and <18.5% for other induced field quantities (local induced E-field,  $2\times2\times2\,\text{mm}^3$  cube averaged induced E-field, 5 mm line averaged induced E-field,  $1\,\text{cm}^2$  area averaged induced current density).

	DASY8/6 Uncerta	inty Budg to IEC/IEEE		sSA	R1g	
Item	Error Description	Unc. Value	Probab.	Div.	$(c_i)$	Std. Unc.
		(±dB)	Distr.			(±dB)
Meas	surement system					
1	Amplitude calibration uncertainty	0.35	N	1	1	0.35
2	Probe anisotropy	0.60	R	$\sqrt{3}$	1	0.35
3	Probe dynamic linearity	0.20	R	$\sqrt{3}$	1	0.12
4	Probe frequency domain response	0.30	R	$\sqrt{3}$	1	0.17
5	Probe frequency linear interp. fit	0.15	R	$\sqrt{3}$	1	0.09
6	Spatial averaging	0.10	R	√3	1	0.06
7	Parasitic E-field sensitivity	0.10	R	$\sqrt{3}$	1	0.06
8	Detection limit	0.15	R	$\sqrt{3}$	1	0.09
9	Readout electronics	0	N	1	1	0
10	Probe positioning	0.19	N	1	1	0.19
11	Repeatability	0.10	N	1	1	0.10
12	Surface field reconstruction	0.20	N	1	1	0.20
Num	erical simulations					
13	Grid resolution	0.02	R	$\sqrt{3}$	1	0.01
14	Tissue parameters	0	R	√3	1	0
15	Exposure position	0	R	$\sqrt{3}$	1	0
16	Source representation	0.09	N	1	1	0.09
17	Convergence and power budget	0	R	$\sqrt{3}$	1	0
18	Boundary conditions	0.10	R	$\sqrt{3}$	1	0.06
19	Phantom loading/backscattering	0.10	R	$\sqrt{3}$	1	0.06
Coml	bined uncertainty $(k=1)$					0.63
Expa	nded uncertainty $(k=2)$					1.27 (33.9%)

Table 1.2: Uncertainty budget for peak 1 g mass-average SAR measured with DASY8/6 Module WPT, assessed according to IEC/IEEE 63184.

	DASY8/6 Uncertai	nty Budgo to IEC/IEEE		sSAI	R10g	l
Item	Error Description	Unc. Value	Probab.	Div.	$(c_i)$	Std. Unc.
		(±dB)	Distr.			(±dB)
Meas	urement system					
1	Amplitude calibration uncertainty	0.35	N	1	1	0.35
2	Probe anisotropy	0.60	R	$\sqrt{3}$	1	0.35
3	Probe dynamic linearity	0.20	R	$\sqrt{3}$	1	0.12
4	Probe frequency domain response	0.30	R	√3	1	0.17
5	Probe frequency linear interp. fit	0.15	R	$\sqrt{3}$	1	0.09
6	Spatial averaging	0.10	R	$\sqrt{3}$	1	0.06
7	Parasitic E-field sensitivity	0.10	R	$\sqrt{3}$	1	0.06
8	Detection limit	0.15	R	$\sqrt{3}$	1	0.09
9	Readout electronics	0	N	1	1	0
10	Probe positioning	0.19	N	1	1	0.19
11	Repeatability	0.10	N	1	1	0.10
12	Surface field reconstruction	0.20	N	1	1	0.20
Nume	erical simulations					
13	Grid resolution	0	R	$\sqrt{3}$	1	0
14	Tissue parameters	0	R	$\sqrt{3}$	1	0
15	Exposure position	0	R	$\sqrt{3}$	1	0
16	Source representation	0.04	N	1	1	0.04
17	Convergence and power budget	0	R	$\sqrt{3}$	1	0
18	Boundary conditions	0.10	R	$\sqrt{3}$	1	0.06
19	Phantom loading/backscattering	0.10	R	$\sqrt{3}$	1	0.06
Comb	pined uncertainty $(k=1)$					0.63
Expai	nded uncertainty $(k=2)$					1.25 (33.4%)

Table 1.3: Uncertainty budget for peak  $10\,g$  mass-average SAR measured with DASY8/6 Module WPT, assessed according to IEC/IEEE 63184.

#### 5.3 Assessment of Peak Incident Fields

#### 5.3.1 Measurement Procedure

The workflow to demonstrate compliance with RL is illustrated in Figure 1.10. Detailed descriptions of each step can be found in Section 7 of the DASY8/6 Module WPT V2.6+ Manual [10]. It is recommended to perform a system check before any compliance testing with the V-Coil source that operates at the frequencies closest to that of the DUT. This provides the confidence that the system operates within its specifications.

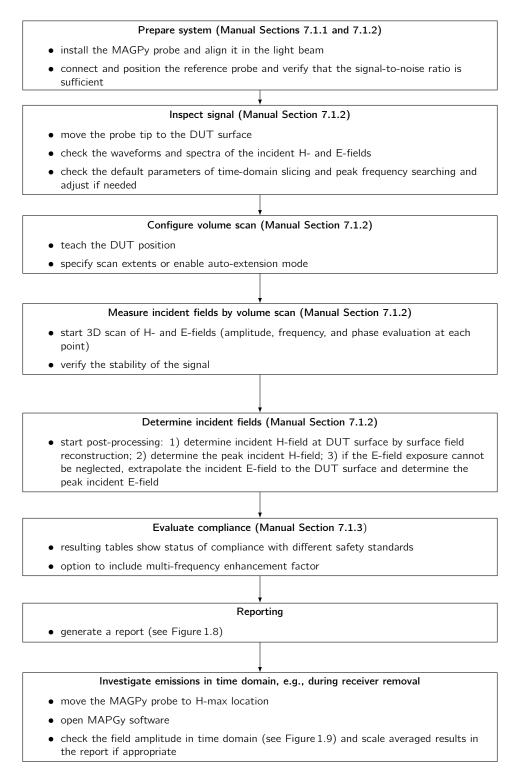


Figure 1.10: Step-by-step measurement procedure for using DASY8/6 Module WPT to evaluate compliance of WPT devices and systems with RL.

#### 5.3.2 Uncertainty

The uncertainty for evaluations performed with DASY8/6 Module WPT V2.6+ was determined according to IEC/IEEE 63184 [1] and documented in the DASY8/6 Module WPT V2.6+ Manual [10]. Typically, the uncertainties (k = 2) are <16.6% for the incident H-field (at the lowest plane of the measurement volume, see Table 1.4) and <24.4% for the incident E-field (see Table 1.5).

	DASY8/6 Uncertainty B according	udget for to IEC/IEEE		ıcide	nt H	l-Field
Item	Error Description	Unc. Value (±dB)	Probab. Distr.	Div.	$(c_i)$	Std. Unc. (±dB)
Meas	surement system	(±ub)	DISTI.			(±ub)
1	Amplitude calibration uncertainty	0.35	N	1	1	0.35
2	Probe anisotropy	0.60	R	$\sqrt{3}$	1	0.35
3	Probe dynamic linearity	0.20	R	√3	1	0.12
4	Probe frequency domain response	0.30	R	$\sqrt{3}$	1	0.17
5	Probe frequency linear interp. fit	0.15	R	$\sqrt{3}$	1	0.09
6	Spatial averaging	0.10	R	$\sqrt{3}$	1	0.06
7	Parasitic E-field sensitivity	0.10	R	$\sqrt{3}$	1	0.06
8	Detection limit	0.15	R	$\sqrt{3}$	1	0.09
9	Readout electronics	0	N	1	1	0
10	Probe positioning	0.19	N	1	1	0.19
11	Repeatability	0.10	N	1	1	0.10
12	Surface field reconstruction	0.30	N	1	1	0.30
Comb	pined uncertainty $(k=1)$					0.67
Expai	nded uncertainty $(k=2)$					1.33 (16.6%)

Table 1.4: Uncertainty budget for peak incident H-field measured with DASY8/6 Module WPT, assessed according to IEC/IEEE 63184.

	DASY8/6 Uncertainty according	Budget to IEC/IEEE		dent	E-Fi	ield
Item	Error Description	Unc. Value	Probab.	Div.	$(c_i)$	Std. Unc.
		(±dB)	Distr.			(±dB)
Meas	urement system				•	
1	Amplitude calibration uncertainty	0.53	N	1	1	0.53
2	Probe anisotropy	0.80	R	√3	1	0.46
3	Probe dynamic linearity	1.00	R	$\sqrt{3}$	1	0.58
4	Probe frequency domain response	0.30	R	$\sqrt{3}$	1	0.17
5	Probe frequency linear interp. fit	0.15	R	$\sqrt{3}$	1	0.09
6	Parasitic H-field sensitivity	0.20	R	$\sqrt{3}$	1	0.12
7	Detection limit	0.15	R	$\sqrt{3}$	1	0.09
8	Readout electronics	0	N	1	1	0
9	Repeatability	0.10	N	1	1	0.10
Comb	bined uncertainty $(k=1)$					0.95
Expa	nded uncertainty $(k=2)$					1.89 (24.4%)

Table 1.5: Uncertainty budget for incident E-field measured with DASY8/6 Module WPT with linear gradients across the probe, assessed according to [1]4.

6. CONCLUSIONS Application Note

#### 6 Conclusions

This application note provides guidance on how to use DASY8/6 Module SAR 16.4+ and DASY8/6 Module WPT V2.6+ for measurement-based compliance testing against BR (induced E-field for nerve stimulation, SAR for thermal effect) in accordance with [1] for inductive WPT devices operating in the frequency range from  $3\,\mathrm{kHz}$  to  $30\,\mathrm{MHz}$ .

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# A Reconstruct a Vector Potential $\vec{A}$ From a Magnetic Field

The induced current simulation requires a vector potential instead of the H-field as input. The vector potential  $\vec{A}$  has the property of curl $(\vec{A}) = \vec{B}$  and, in free space, of  $\vec{B} = \mu_0 \vec{H}$ .

This equation is non-trivial, given that any additional gradient field fulfills  $\operatorname{curl}(\vec{A} + \operatorname{grad}(\phi)) = \vec{B}$ . Luckily, an explicit formula that requires the use of  $\operatorname{div}(\vec{B}) \equiv 0$  on a rectilinear grid can be derived [14]. For the x-component (others are cyclic permutations), the formula reads:

$$A_{x} = -\int_{0}^{y} \left[ \frac{1}{3} B_{z}(x, y, z) + \frac{1}{6} B_{z}(x, y, 0) \right] dy + \int_{0}^{z} \left[ \frac{1}{3} B_{y}(x, y, w) + \frac{1}{6} B_{y}(x, 0, w) \right] dw$$
 (2)

where the subscripts x, y, and z denote the components along their corresponding axes, and v and w denote the integration variables. The axes origin, i.e., the point where the path integrals begin to integrate, can be arbitrarily chosen. Currently, the most dominant location of the  $\vec{B}$  field is chosen as the origin to minimize numerical integration artefacts. Details are provided in [14].

#### B Effect of Backscattering on the Source

#### **B.1** Objectives

In this section, we assess the effect of the phantom loading or backscattering on the incident field for frequencies below 4 MHz by comparing the dissipated energy in the phantom to the H-field energy. The uncertainty in the determination of the induced fields due to the incident field without the phantom is also determined.

#### B.2 Theory

Maxwell's equation in the frequency domain (with linear constitutive material models) reads (see Sim4Life manual for details):

$$\nabla \times \vec{E} = -j\omega \vec{B} = -j\omega \mu \vec{H} \tag{3a}$$

$$\nabla \times \vec{H} = -j\omega \vec{D} + \vec{J} = -j\omega \epsilon \vec{E} + \sigma \vec{E} + \vec{J_0}$$
(3b)

$$\nabla \times \vec{D} = \nabla \cdot \epsilon \vec{E} = \rho \tag{3c}$$

$$\nabla \times \vec{B} = \nabla \cdot \mu \vec{H} = 0 \tag{3d}$$

With a vector potential  $\vec{A}$  defined as  $\nabla \times \vec{A} = \vec{B} = \mu \vec{H}$  (in the Coulomb gauge, i.e.,  $\nabla \cdot \vec{A} = 0$ ), the E-field can be written as  $\vec{E} = -j\omega \vec{A} - \nabla \phi$ , where  $\phi$  is an additional scalar potential. The complex permittivity  $\tilde{\epsilon} := \epsilon + \frac{\sigma}{j\omega}$  and the divergence-freeness of the  $\vec{A}$  allows the  $\nabla \times \vec{H}$  equation to be rewritten as

$$\nabla \times \frac{1}{\mu} \nabla \times \vec{A} = \underbrace{\omega^2 \tilde{\epsilon} \vec{A} - j\omega \tilde{\epsilon} \nabla \phi}_{:=\omega - \text{terms}} + \vec{J_0}$$
(4)

The H-field  $\vec{H}$  is the static H-field, i.e., it is not altered by the induced E-field, if

$$\nabla \times \frac{1}{u} \nabla \times \vec{A} = \vec{J_0},\tag{5}$$

i.e., the two  $\omega$ -terms are negligible. In the following, the order-of-magnitude scalings of those 2  $\omega$ -terms are investigated. The order of magnitude can be estimated by means of the in-order-of notation  $\mathcal{O}(\cdot)$ .

Since in a vacuum there are no free charges, i.e.,  $\rho$  vanishes, the scalar and vector potential are related as  $\nabla \times \epsilon \nabla \phi = -j\omega \nabla \cdot \epsilon \vec{A}$ . Given a characteristic length scale L to estimate the spatial derivations, the following relationship is provided:  $\mathcal{O}(\tilde{\epsilon}\phi/L^2) = \mathcal{O}(\omega\tilde{\epsilon}A/L)$ , i.e,  $\phi$  scales like  $\phi = \mathcal{O}(\omega AL)$ . Application of the same scaling strategy to two  $\omega$ -terms in (4) yields  $\mathcal{O}(\omega^2\tilde{\epsilon}A)$  in both cases. Therefore, it can be estimated

$$\frac{\text{both-}\omega\text{-terms}}{\nabla \times \frac{1}{\mu} \times A\text{-term}} = \mathcal{O}(\omega^2 \tilde{\epsilon} \mu L^2), \tag{6}$$

i.e., written with permittivity and conductivity, the  $\omega$ -terms can be neglected when

$$\omega^2 \epsilon \mu L^2 \ll 1 \tag{7a}$$

$$\omega \sigma \mu L^2 \ll 1 \tag{7b}$$

Calculated values for the two  $\omega$ -terms of the tissue material properties  $\epsilon_r = 55$ ,  $\sigma = 0.75\,\mathrm{S/m}$  and a reference coil diameter, where the coil diameter was used as the characteristic length, are provided in Table 6. At 4 MHz, both values are much smaller than  $-20\,\mathrm{dB}$ , i.e., the quasi-static conditions can still be considered as valid.

The fist criterion (7a) can be rewritten using the wave-length  $\lambda$ , the frequency  $f=\omega/2\pi$ , the speed of light within the phantom c and the relations  $\sqrt{1/(\epsilon\mu)}=c=\lambda\cdot f=\lambda\cdot\omega/(2\pi)$ , i.e., replacing  $\omega^2\epsilon\mu$  with  $(2\pi/\lambda)^2$ 

$$\frac{2\pi L}{\lambda} \ll 1 \quad \Longleftrightarrow \quad L \ll \frac{\lambda}{2\pi} \tag{8}$$

The second criterion in (7b) can be further simplified using the skin depth  $\delta = \sqrt{2/(\omega \sigma \mu)}$  (valid if  $\omega \ll \sigma/\epsilon$ ), i.e., replacing  $\omega \sigma \mu$  with  $2/\delta^2$ :

$$\frac{\sqrt{2}L}{\delta} \ll 1 \quad \Longleftrightarrow \quad L \ll \frac{\delta}{\sqrt{2}} \tag{9}$$

SPEAG, DASY8/6 Application Note: Testing Compliance with IEC/IEEE 63184, May 2024

Frequency	Coil diameter	$\omega^2 \epsilon \mu L^2$	$\omega \sigma \mu L^2$
[kHz]	[mm]	[dB]	[dB]
3	454.0	-147.0	-48.7
85	200.0	-103.1	-33.9
400	52.5	-99.5	-43.7
1000	52.5	-83.5	-35.7
2000	52.5	-71.5	-29.7
4000	52.5	-59.5	-23.7

Table 6: Results of the calculations of the two  $\omega$ -terms (i.e.,  $\omega^2 \epsilon \mu L^2$  and  $\omega \sigma \mu L^2$ ) for tissue material properties  $\epsilon_r = 55$ ,  $\sigma = 0.75 \, \text{S/m}$  and a coil diameter of 52 mm. The coil diameter was used as the characteristic length.

#### **B.3** Simulation Evidence

As a next step, we simulate the extent of back-scattering or the loading by the phantom by comparing the energy absorbed in the phantom to the maximum stored energy in the H-field. This ratio is expressed as

$$Q^{-1} = \frac{\text{power absorpbed in the phantom}}{2\pi f(\text{maximum energy stored})} = \frac{\int_{V} (\rho \text{ SAR}) dV}{2\pi f W_{H}}$$
 (10)

In Eqn. (10), the magnetic energy stored  $W_{\rm H}$  is calculated by integrating the product of the H-field strength and the magnetic flux density over a volume that is sufficiently large for convergence, and the absorbed power in the phantom is calculated by integrating the product of the mass density and SAR over the entire volume of the phantom.

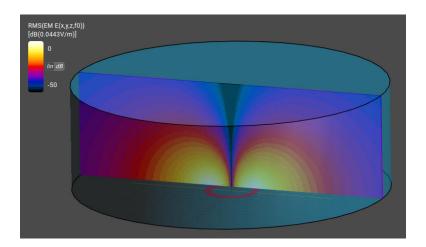


Figure 11: Normalized E-field of a generic WPT transmit coil with a diameter of 100 mm induced in a cylindrical phantom with a conductivity of 0.75 S/m placed at a distance of 4 mm.

Figure 11 shows an example of the E-field distribution in a cylindrical phantom positioned at a distance of 4 mm above a generic coil with a diameter of 100 mm simulated with the magneto quasistatic solver. According to the current draft of IEC 63184 [1], the relative permittivity of the tissue simulating liquid (TSL) in the phantom is  $\epsilon_r = 55$ , the conductivity  $\sigma = 0.75\,\text{S/m}$ , and the density  $\rho = 1000\,\text{kg/m}^3$ . The effect of the phantom on the total incident H-field along the vertical center line of the coil is shown in Figure 12. The  $Q^{-1}$  values according to Eqn. (10) for the generic coil at frequencies of 400 kHz and 6.78 MHz are given in Table 7.

Frequency	$W_{H}$	$\int_{V} (\rho  SAR) dV$	$Q^{-1}$
[kHz]	$[\mu J]$	[mW]	[dB]
400	2.1	0.84	-30
6780	2.1	240	-18

Table 7: Maximum H-field energy and dissipated power in the phantom per  $1A_{peak}$  for the 400 kHz and 6.78 kHz verification sources. The load of the phantom is <1% at frequencies <4 MHz.

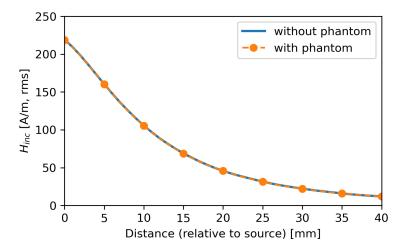
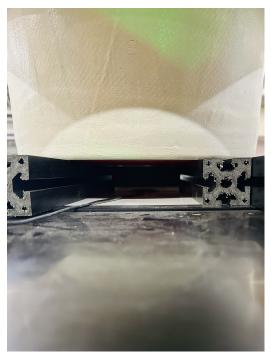


Figure 12: Comparison of the total incident H-fields (per  $1A_{peak}$ ) along the z-axis center line of the coil with and without the phantom for the reference source V-Coil 50/400.



(a) Perspective view showing the probe measuring in  $\ensuremath{\mathsf{TSL}}$ 

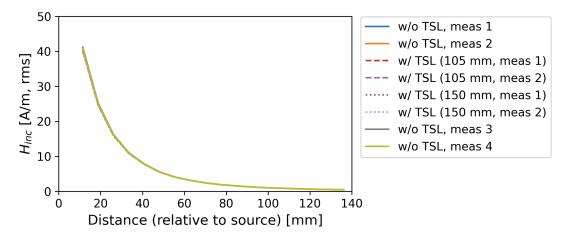


(b) Side view showing the placement of the phantom and the source

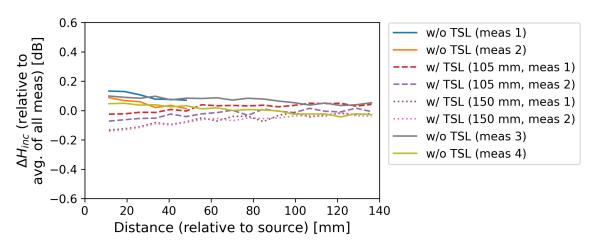
Figure 13: Setup for experimental confirmation of incident H-field insensitivity to the presence of the phantom.

#### **B.4** Experimental Confirmation

To experimentally confirm that the effect of the phantom on the incident H-field is very small, measurements were made with a specially sealed MAGPy probe placed inside the ELI phantom filled with tissue simulating liquid (TSL) HBBL4-250V3, which has the nominal values of  $\epsilon_r = 55$  and  $\sigma = 0.75\,\mathrm{S/m}$ . The phantom was placed above the 400 kHz reference source (i.e., V-Coil50/400) at d  $\approx 2\,\mathrm{mm}$ . The current fed to the coil was determined by measuring the voltage across the current monitoring resistor with an oscilloscope. A photo of the setup is shown in Figure 13. The probe was moved to different distances (with a DASY robot) along a vertical observation line. The incident H-field was measured at each distance for three cases: (1) without TSL, (2) with TSL of a filling depth of 105 mm, and (3) with TSL of a filling depth of 150 mm. The measurement for the case without TSL was repeated four times. The measurements for the cases with TSL were repeated twice, and were made between the second and third measurements for the case without TSL. The H-field measurement results are illustrated in Figure 14. The deviations are well within the expected measurement repeatability of <0.5 dB and confirms the theoretical considerations that the effect is less than 1% or 0.1 dB.



(a) Total H-fields along the vertical line for three TSL-filling cases



(b) Deviations in the total H-field along the vertical line for three TSL-filling cases

Figure 14: Measurement results of the incident H-field, confirming its insensitivity to the presence of the phantom. The results showed include three cases: with TSL of a filling depth of 105 mm, with TSL of a filling depth of 150 mm, without TSL.

#### **B.5** Conclusions

The effect of the phantom loading or backscattering is less than 1% for frequencies <4 MHz as derived from theoretical considerations, supported by simulations and verified by measurements. Therefore, when the evaluation is performed on the incident field only, i.e., without phantom, the coupling between phantom and source results in an additional uncertainty of  $0.1\,\mathrm{dB}$ . Neglecting the flux cancellation due to the induced current in the phantom leads to a small overestimation and therefore is conservative and is not considered in the uncertainty estimation.

#### C Validation of Total Field Evaluation of DASY8/6 Module WPT

#### C.1 Evaluation of the Induced Fields by the Validation Source V-Coil50/6780 V2

#### C.1.1 Instrumentation and V&V Sources

The configuration of the DASY8/6 Module WPT system used in the validation measurements is listed in Table 8.

System	Type:	DASY6 Module WPT
	Software Version:	V2.6
	Manufacturer:	Schmid & Partner Engineering AG, Switzerland
Positioner	Robot:	TX90 XL
	Serial No:	F/18/0004593/A/001
	Controller:	CS8C
	Serial No:	F/18/0004593/C/001
	Manufacturer:	Stäubli, France
Probe	Type:	MAGPy-8H3D+E3D V2
	Serial Number:	3065
	Calibrated On:	Apr. 6, 2023
	Next Calibration:	Apr. 2024
	Frequency Range:	3 kHz–10 MHz
	H-Field Dynamic Range:	0.1-3200 A/m
	E-Field Dynamic Range:	0.1–2000 V/m
	H-Field Sensor Area:	1 cm <sup>2</sup>
	E-Field Sensor Length:	5 cm
	Probe Length:	335 mm
	Probe Tip Diameter:	60 mm (flat tip)
	Manufacturer:	Schmid & Partner Engineering AG, Switzerland
6.78 MHz Verification Source	Source Model:	V-Coil50/6780 V2
	Source Serial No.:	1014
	Source Dimensions:	$250\mathrm{mm} imes125\mathrm{mm} imes35\mathrm{mm}$
	Source Output Freq.:	6.78 MHz
	Source Current:	0.394 A
	Source Evaluated On:	Jan. 29, 2024
	Source Manufacturer:	Schmid & Partner Engineering AG, Switzerland

Table 8: DASY6 Module WPT system and Validation Source

#### C.1.2 Method

The 6.78 MHz validation source was simulated with the fullwave finite-difference time-domain solver and the magneto quasi-static (MQS) solver in Sim4Life V7.2, and also measured using DASY6 Module WPT V2.6. The total field approximation was also applied with the incident E-field obtained from simulation/measurement as the input. This is a nearly worst-case evaluation, as the contributions of the incident E-field to the induced E-field is further reduced at lower frequencies.

#### C.1.3 Results

The results are summarized in Table 9.

Simulation vs	Method	pE <sub>ind</sub> a	psSAR <sub>1g</sub>	psSAR <sub>10g</sub>
Measurement		[V/m]	[W/kg]	[W/kg]
Simulation	Fullwave	109	4.05	1.97
	MQS	107	4.06	1.95
	Total field approximation	107	4.06	1.95
Measurement	MQS	104	3.97	1.97
	Total field approximation	104	3.97	1.97

<sup>&</sup>lt;sup>a</sup> Maximum induced E-field

Table 9: Results of the induced field evaluations performed with Sim4Life V7.2 and DASY6 Module WPT V2.6.

#### C.2 Conclusions

The implemented MQS assessment provides accurate results for the fields induced by the incident H-field only. The total field approximation implemented in DASY8/6 Module WPT V2.6+ provides a reliable assessment of the maximum induced fields, e.g., psSAR1g/10g and  $pE_{ind}$ .

D Verification Report of Low Frequency Magneto Quasi-Static Solver



### Verification Report MQS001AA201507

# Sim4Life and SEMCAD X Low Frequency Magneto Quasi-Static Solver

George Tsanidis<sup>1</sup>, Theodoros Samaras<sup>2</sup>

#### Thessaloniki, July 2015

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#### **Executive Summary**

The THESS S.A. was mandated by ZMT Zurich MedTech AG (Offer No. 1420) to independently verify the Sim4Life and SEMCAD X platforms. The here documented MQS001AA201507 verification benchmark for the EM LF Magneto Quasi-static Solver was developed and tested for the Sim4Life Version v2.0 and SEMCAD X v14.8, and documented such that ZMT can automatically run the test for any new software version.

The MQS001AA201504 benchmark verifies the EM LF Magneto Quasi-static Solver against the analytically calculated value of the magnetic field generated by a circular thin-wire coil inside an adjacent conducting sphere. This benchmark tests the following solver features:

- that the EM LF Magneto Quasi-static Solver converges to the right solution
- that material interfaces are properly handled
- that the Biot-Savart sources are correctly implemented
- that post-processing correctly calculates derived quantities such as current density.

The agreement of the theoretically calculated magnetic field with the values derived from the Sim4Life and SEMCAD X platform is very good: The current density values at different positions inside the sphere match for the two methods with a deviation smaller than 0.5% for the finer discretization of the computational domain. With increasing grid resolution, the simulation results converge to the analytical solution.

In conclusion, the numerical MQS solver of Sim4Life and SEMCAD X therefore meets the requirements for modeling the magneto-quasistatic equation.

# Contents

1	Objectives
2	Methodology
	2.1 Introduction
	2.2 Analytical solution
	2.3 Numerical Modeling
3	Results
	3.1 Criterion of convergence
	3.2 Grid step
	3.3 Material interfaces
1	Conclusion

#### 1 Objectives

The objective of this verification report MQS001AA201507 is to document the verification of the Sim4Life v2.0 and SEMCAD X v14.8 Low Frequency Magneto Quasi-static Solver by comparing numerical to analytical solutions of a specific problem.

The MQS solver first calculates a magneto-static vector potential  $(A_0)$  using the Biot-Savart law and subsequently determines the induced E-fields and currents using potential continuity while considering the inhomogeneous dielectric property distributions in the human anatomy [2]. The equation  $\nabla \cdot \sigma \nabla \phi = -j\omega\nabla \cdot (\sigma A_0)$  is solved  $(\sigma$ : conductivity,  $\omega$ : angular frequency,  $\phi$ : electric scalar potential) which is valid at frequencies where ohmic currents dominate over displacement currents.

The following features of the EM LF Magneto Quasi-static solver have been identified as fundamental and requiring verification:

- that the EM LF Magneto Quasi-static Solver converges to the right solution
- that material interfaces are properly handled
- that the Biot-Savart sources are correctly implemented
- that post-processing correctly calculates derived quantities such as current density

For that purpose an analytically solvable benchmark case has been chosen such that it makes use and covers all of these critical features. The benchmark is a homogeneous sphere exposed to current carrying ring wire.

#### 2 Methodology

#### 2.1 Introduction

The field of Magnetostatics was widely studied during the 19th century. The work of J.B. Biot and F. Savart made possible the calculation of the magnetic field originating from an electric current. They provide an approximation of the Maxwell equations valid for low frequency, provided a quasi-static behavior that can be assumed in the case of slow time variations (low frequency) and sufficiently small dimensions. The approximation condition is  $\left(\frac{d}{\lambda}\right)^2 << 1$ , where d is the diameter of the computational domain and  $\lambda$  the wavelength. The law of Biot-Savart can be used in order to easily calculate the value of the magnetic field inside a head-sized sphere, with the dielectric properties of the human brain, created by an adjacent circular loop coil.

The numerically derived results of the EM LF Magneto Quasi-static Solver of the Sim4Life and SEM-CAD X platform are compared with the theoretically calculated results in order to evaluate the reliability and accuracy of the former and this deviation is examined with respect to

- the relative solver tolerance used to terminate the numerical process,
- the grid step of the computational domain,
- the implementation of the material interfaces

#### 2.2 Analytical solution

A surface coil (loop of uniform current) adjacent to a homogeneous conducting sphere can be used in order to predict the performance of MRI surface coils close to the human head. By solving the inhomogeneous boundary value problem of the system, the electromagnetic field inside the sphere can be calculated. The sphere's parameters, i.e. the relative dielectric constant  $\epsilon_r$  and the conductivity  $\sigma$ , are chosen so as to model the human brain.

The magnetic field produced by the surface coil adjacent to the homogeneous sphere has been calculated by solving the inhomogeneous boundary value problem of a ring of radius R carrying uniform current I adjacent to a conducting dielectric sphere of radius  $\alpha$  centered at the origin of a spherical coordinate system (Fig. 1).

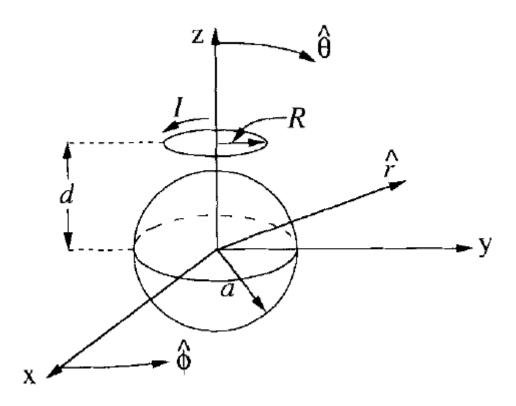


Figure 1: Schematic arrangement of a homogeneous conductive sphere near a circular current loop. The sphere is centered at the origin of the spherical coordinate system

The magnetic field inside the sphere is given by [1]

$$B_r(r,\theta,\phi) = \sum_{i=1}^{\infty} b_{l0} \sqrt{\frac{l(l+1)(2l+1)}{4\pi}} P_i(\cos\theta) \frac{j_i(k^{in}r)}{k^{in}r}$$
(1)

$$B_{\theta}(r,\theta,\phi) = -\sum_{i=1}^{\infty} b_{l0} \sqrt{\frac{(2l+1)}{l(l+1)4\pi}} sin\theta \frac{dP_i(cos\theta)}{dcos\theta} \frac{1}{k^{in}r} \frac{\vartheta(j_i(k^{in}r))}{\vartheta r}$$
(2)

$$B_{\phi}(r,\theta,\phi) = 0 \tag{3}$$

where  $j_i(kr)$  denotes the spherical Bessel functions of the first kind and  $P_i(\cos\theta)$  are Legendre polynomials.

The wavenumber inside the sphere  $k^{in}$  is obtained from Maxwell's equation ,

$$\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon \frac{\partial \vec{E}}{\partial t}.$$
 (4)

where  $\mu_0$  is the permeability of free space. Substituting  $\vec{J} = \sigma \vec{E}$  and defining the relative dielectric constant as  $\epsilon_r = \frac{\epsilon}{\epsilon_0}$  equation (4) is transformed to

$$i\omega\nabla \times \vec{B} = (i\omega\mu_0\sigma + \epsilon_r \frac{\omega^2}{c^2})\vec{E}$$
 (5)

where  $\omega$  is angular frequency,  $\sigma$  is electrical conductivity and c is the speed of light. The wave number inside the sphere  $k^{in}$  is the square root of the coefficient of

$$(k^{in})^2 = i\omega\mu_0\sigma + \epsilon_r \frac{\omega^2}{c^2} \tag{6}$$

Solving equation (5) for the electric field, one is allowed to express the current density  $\vec{J}$  as

$$J_{\phi}(r,\theta,\phi) = \frac{i\omega\sigma}{k^{in}} \sum_{i=1}^{\infty} b_{l0} \sqrt{\frac{(2l+1)}{l(l+1)4\pi}} j_i(k^{in}r) sin\theta \frac{dP_i(cos\theta)}{dcos\theta}$$
 (7)

The  $b_{l0}$  are found by satisfying the boundary conditions of the magnetic field at the surface of a homogeneous conducting, dielectric sphere given an incident magnetic field produced by an adjacent ring of uniform current

$$b_{l0} = \mu_0 I 2\pi \sqrt{\frac{(2l+1)}{l(l+1)4\pi}} \frac{(k^{out})^2 R^2 h_i^{(1)} (k^{out} \sqrt{d^2 + R^2})}{\sqrt{d^2 + b^2}} \frac{dP_i(\xi)}{d\xi}$$

$$\frac{k^{in} (j_i (k^{out} \alpha) y_{i+1} (k^{out} \alpha) - y_i (k^{out} \alpha) j_{i+1} (k^{out} \alpha))}{k^{in} h_i^{(1)} (k^{out} \alpha) j_{i+1} (k^{in} \alpha) - k^{out} j_i (k^{in} \alpha) h_{i+1}^{(1)} (k^{out} \alpha)}$$
(8)

where  $k^{out}$  is the wave number in free space,  $y_i$  and  $h_i^{(1)}$  are spherical Bessel functions of the second and third kinds, respectively, and  $\xi$  is the cosine of the angle subtended by the loop  $\xi = \frac{d}{\sqrt{d^2 + R^2}}$ 

#### 2.3 Numerical Modeling

This verification study intends to compare the current density J calculated by the analytical solution, with the numerical results obtained by Sim4Life and SEMCAD X for the same problem. A sphere of 60mm radius and a surface coil of 20mm radius were placed at a distance of 50mm between the coil center and the nearest point of the sphere. For mathematical convenience the z axis of the coordinate system was chosen to be parallel to the axis of the coil. The sphere's relative dielectric permittivity  $\epsilon_r$  and electrical conductivity  $\sigma$  are that of a human brain (white matter:  $\sigma = 0.0626 \, S/m$ ,  $\epsilon_r = 69800$ ).

#### 3 Results

#### 3.1 Criterion of convergence

No difference in the numerical results was observed when changing the value of relative tolerance (which is the criterion of convergence for the computational process) from  $10^{-6}$  to  $10^{-12}$ . In particular the maximum difference between the values of the electric field induced inside the sphere, as extracted at the center, is 0.000172%.

#### 3.2 Grid step

Simulations with uniform grids and varying grid step were performed (0.3, 0.5, 1, 2, 3, 4 and 5 mm) and they gave similar results (Fig. 2). For comparison the magnetic field along two axes (Fig. 3) was extracted and is presented with the theoretical, as is calculated by the Law of Biot-Savart (Fig. 4) and (Fig. 5).

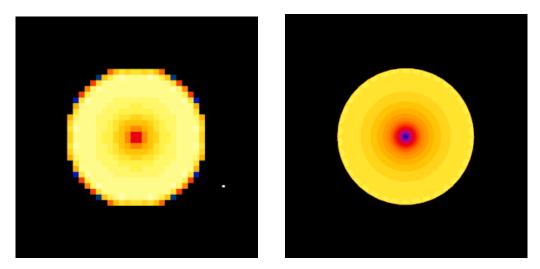


Figure 2: Slice views of the electric field inside the sphere for coarse (Grid Step: 5mm) and fine grid (Grid Step: 0.3mm).

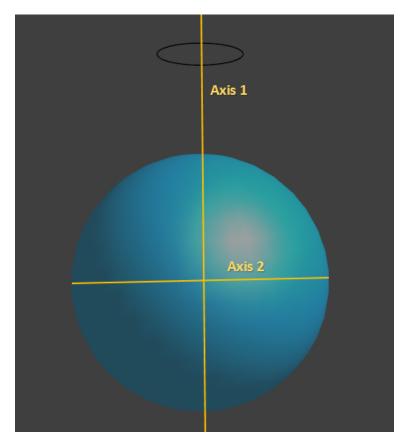


Figure 3: The two axes of extraction.

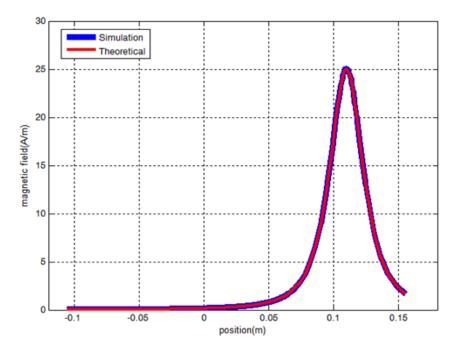


Figure 4: The extracted and the theoretical magnetic field along Axis 1 (Grid step: 0.3mm)

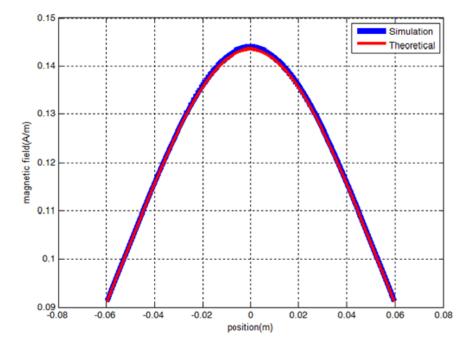


Figure 5: The extracted and the theoretical magnetic field along Axis 2 (Grid step:  $0.3 \mathrm{mm}$ )

For the comparison between the theoretical and the simulated current distribution at the sphere, the deviation of the numerical solution from the analytical solution was evaluated at circles of distances 5, 10, 30 and 50mm from the center of the sphere (Fig. 6).

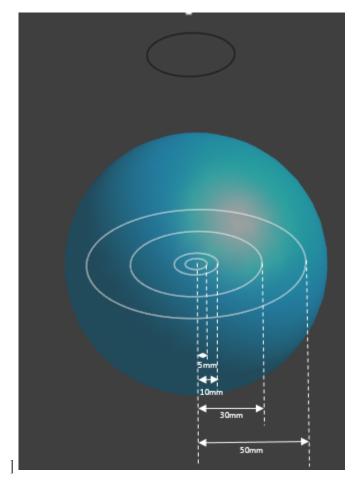


Figure 6: The two axes of extraction.

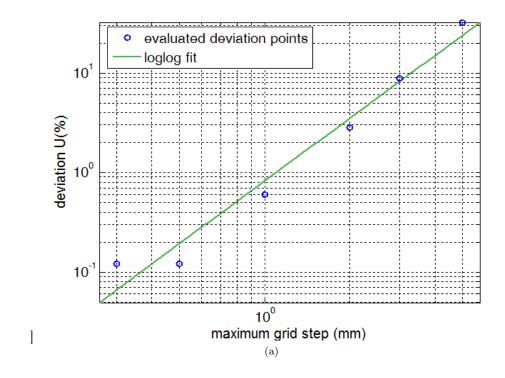
The computational space was discretized with a uniform grid of variable maximum step and the deviation between the numerical (N) and the analytical (T) solution for the current density was evaluated at the points of the numerical solution:

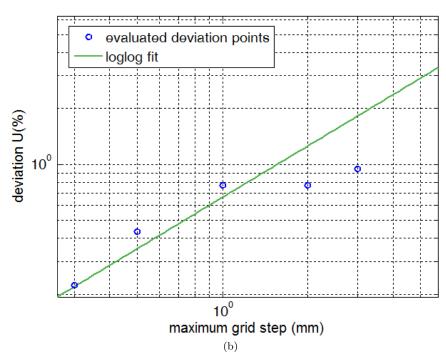
$$deviation U = |\frac{N-T}{T}|100\%$$
(9)

The maximum deviation on the above circles are shown in (Fig. 7). It is clear that, in every case, for a finer discretization of the grid (smaller grid step) the deviation continuously decreases and the numerical solution converges to the analytical solution.

#### 3.3 Material interfaces

For the model with uniform grid step of 0.1mm, the radial component of the magnetic field was extracted on either side of the surface of the sphere, along a circle. At material interfaces, i.e., at the interface of air with the sphere, the radial component of the magnetic field should be identical when approaching the interface from both sides. The absolute value of the relative difference between the two values is shown in (Fig. 8). The deviation never exceeds 0.14% indicating that material interfaces are correctly handled by the solver.





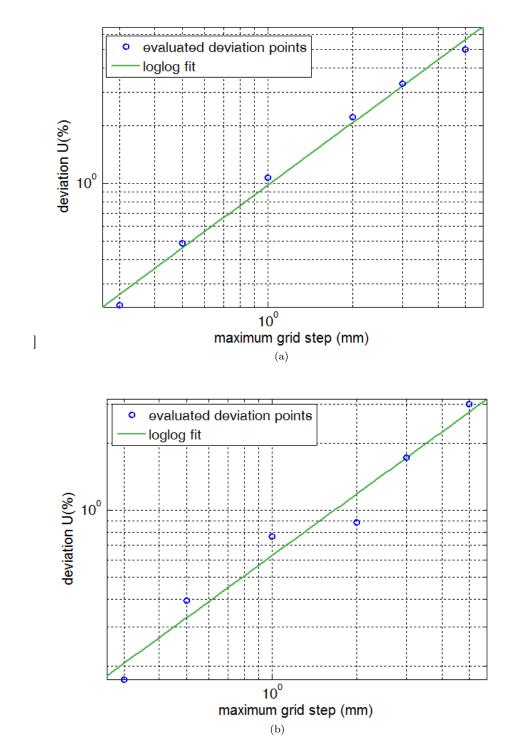


Figure 7: Maximum deviation (see equation (9)) of the numerical from the analytical solution along a circle at distance of 5mm (a), 10mm (b) 30mm (c) and 50mm (d) from the center of the sphere as a function of the maximum grid step.

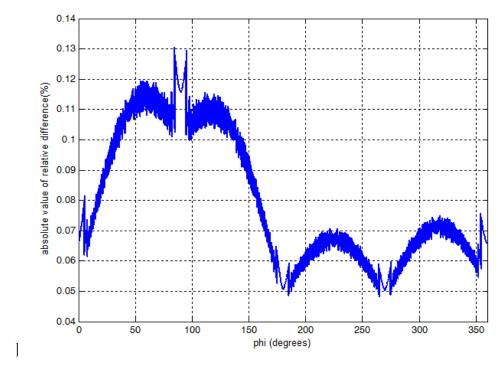


Figure 8: Relative difference (absolute value) between the perpendicular magnetic field components on either side of the sphere's surface. (Grid step: 0.1mm).

#### 4 Conclusion

The purpose of this verification study was to examine the agreement between the numerical results obtained by the Sim4Life and SEMCAD X LF MagnetoQuasi-static solver and the analytically obtained results. It was shown that grid resolution has an important impact on accuracy. It is possible to keep the deviation between numerical and analytical solutions lower than 0.5%, by choosing the appropriate discretization (grid step). With increasing resolution, the simulation results converge to the analytical solution. Proper numerical convergence has been ascertained by varying the convergence criterium. The benchmark case tests the following fundamental solver features:

- that the EM LF Magneto Quasi-static Solver converges to the right solution
- that material interfaces are properly handled
- that the Biot-Savart sources are correctly implemented
- that post-processing correctly calculates derived quantities such as current density

The current density values at different positions inside the sphere match for the two methods with a deviation smaller than 0.5% for the finer discretization of the computational domain. With increasing resolution, the simulation results converge to the analytical solution.

Neither the spatial discretization nor the solution algorithm employed by Sim4Life and SEMCAD X use any assumptions based on the shape of the computational domain. This renders the approach suitable for any complex structures which might occur in biomedical applications. Because of this generalized verification approach it is valid to expect similarly accurate performance of the solver in simple geometries like the presented benchmarks and in more complex geometrical models.

In conclusion, the numerical MQS solver of Sim4Life and SEMCAD X therefore meets the requirements for modeling the magneto-quasistatic equation.

#### References

- [1] J.R. Keltner. Electromagnetic Fields of Surface Coil in Vivo NMR at High Frequencies, p.467-480, Magnetic Resonance in Medicine 22, 1990.
- [2] ZMT. Sim4Life User Manual, ZMT Zurich Med Tech, Zurich, Switzerland, 2015.

#### E Procedure to Test Two-Coil WPT Systems

The test procedure to test WPT systems with two coils working at different fundamental frequencies is outlined here.

- set the peak frequency search range to only capture the fundamental frequency of the first coil
- perform a volume scan and determine the induced fields due to the first coil according to the procedure described in Section ??
- set the peak frequency search range to only capture the fundamental frequency of the second coil
- perform a volume scan and determine the induced fields due to the second coil according to the procedure described in Section ??
- combine the exposure ratios for the induced fields due to the two coils